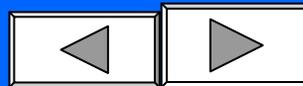


有限元法 电子教案



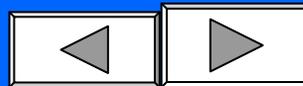
之二



有限元法电子教案

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Tel. 67392760



本课件包括五部分:

一、绪论

第一章 绪论

二、弹性力学基础

第二章 基本概念与假设

第三章 平面问题的基本理论

三、有限元理论及程序

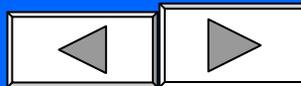
第四章 有限元法的基本概念

第五章 等参元

第六章 三角形单元计算机程序

四、有限元的扩展

五、应用

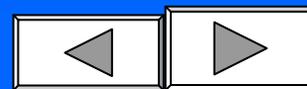


二、弹性力学基础

Foundation of Elasticity Theory

第二章 基本概念与假设

第三章 平面问题的基本理论



参考书

1 简明弹性力学教程 徐芝论 著
高教出版社

2 弹性力学 吴家龙 著
同济大学出版社



第2章 基本概念与假设

弹性力
学:

研究在外界因素（力、热）作用下弹性体的应力，应变和位移。



§ 2.1 基本概念

外力, 应力, 形变, 位移

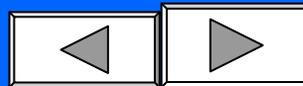
1. 外力 External Force

体积力: 如重力, 惯性力, 用 $F_x, F_y, F_z(x, y, z)$ 表示, 沿坐标方向为正, 因次: $[\text{力}] \cdot [\text{长度}]^{-3}$

Body Force

表面力: 液体压力, 接触力, 用 $T_x, T_y, T_z(x, y, z)$ 表示, 沿坐标方向为正, 因: $[\text{力}] \cdot [\text{长度}]^{-2}$

Surface Force



2.应力

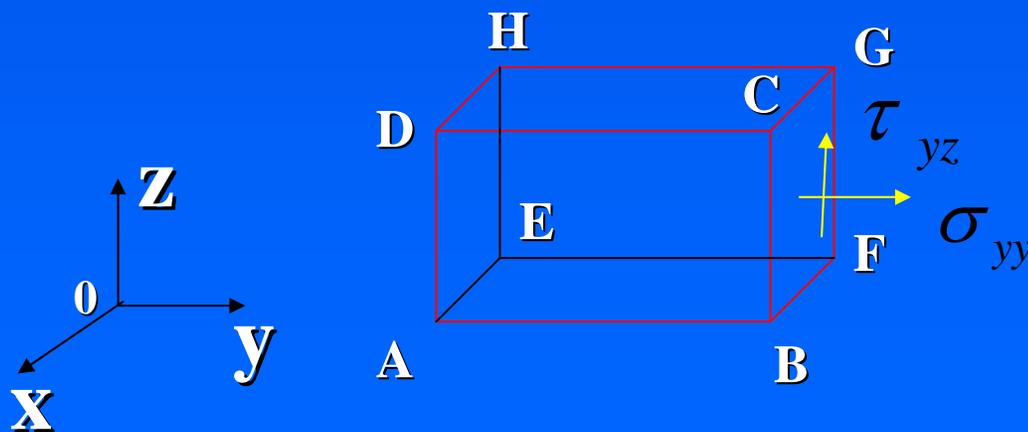
Stress

正应力 σ : $\sigma_x, \sigma_y, \sigma_z$

剪应力 τ : τ_x, τ_y, τ_z (材力)

τ_{ij} : i 表示作用面的法线

j 表示力的方向



σ 的正、负号定义与材料力学一致， τ 则不同

(1)应力所在面的外法线与坐标轴的正方向一致时，

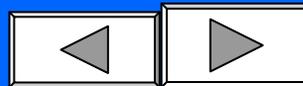
该面为正面，反之为负面。

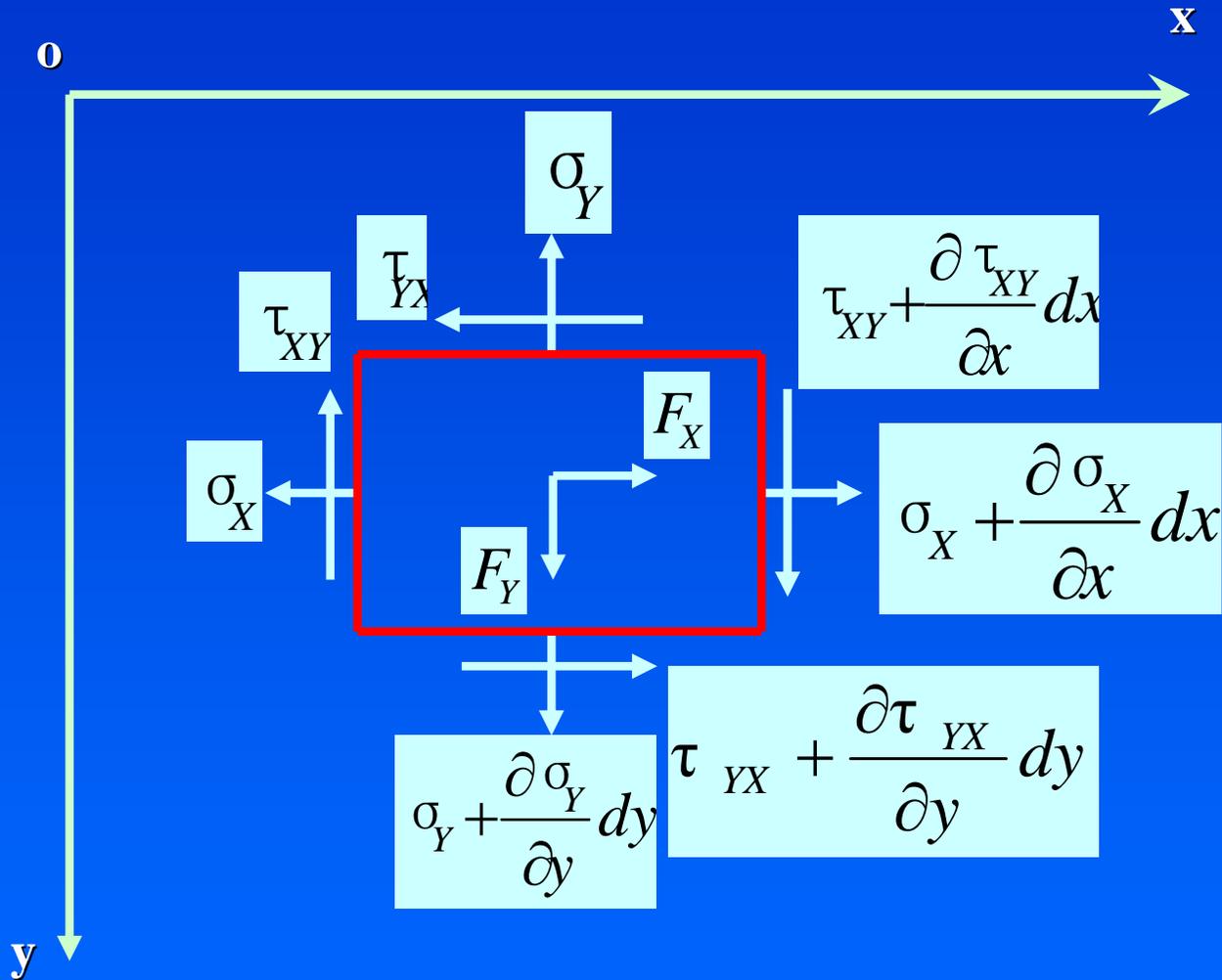
(2)正面上应力的方向与坐标轴的正方向一致时，

应力为正，反之为负。

(3)负面上应力的方向与坐标轴的负方向一致时，

应力为正，反之为负。





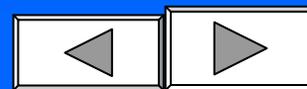
3. 应变 Strain

正应变 ε : $\varepsilon_x, \varepsilon_y, \varepsilon_z$ 伸长为正 (无因次)
剪应变 γ : $\gamma_{xy}, \gamma_{xz}, \gamma_{yz}$ 以直角变小为正

4. 位移 Displacement

移:

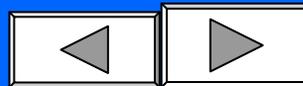
用 u, v, w 表示, 沿坐标正向为正, 反向为负。



§ 2.2 基本假

- (1) 连续体假定 Continuous: 整个物体被介质填满, $\sigma, \varepsilon, \mu = f(x, y, z)$ 可求导, f', f'' 存在。
- (2) 均匀性假定: Homogeneous: 整个物体由同一材料组成, 物体中各处有相同的弹性
- (3) 完全弹性: Perfectly elastic: 服从 HOOKE'S LAW, $\sigma \propto \varepsilon$
- (4) 各向同性: isotropic: 物体内一点的弹性在各个方向都相同。只有两个独立常数 $E, \mu (G)$ 。

符合以上四个假定的物体都称为理想弹性体



(5)小变形假定 Small Deformation

应变与转角都远小于1

- a. 平衡方程可在变形之前的状态上建立.
- b. 所有的代数方程和微分方程都可简化为线性方程
- c. 叠加原理可用



第3章

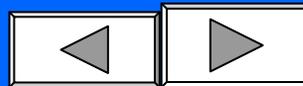
平面问题的基本理论

(Theory of Elasticity for Plane Problems)



主要内容

1. 建立平面应力和平面应变问题的概念;
2. 建立平衡微分方程, 几何方程, 物理方程, 边界条件;
3. 圣维南原理;
4. 求解方法,
按位移和按应力求解平面问题



§ 3.1 平面应力与平面应变问题

(Plane Stress and Plane Strain Problems)

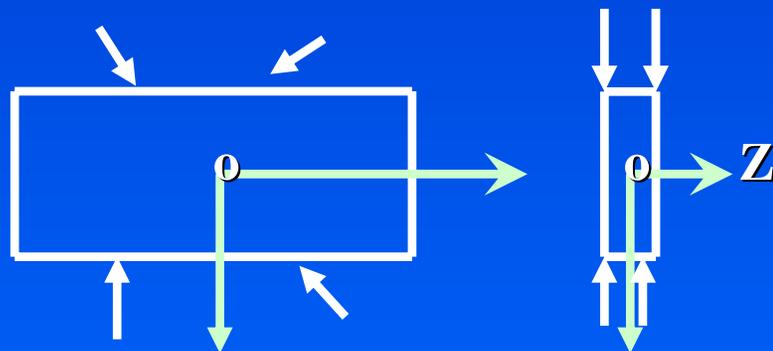
受力体的几何形状和外力在一个方向上保持不变---平面问题

1. 平面应

板等厚度且很薄，外力平行于板面，不沿厚度变化

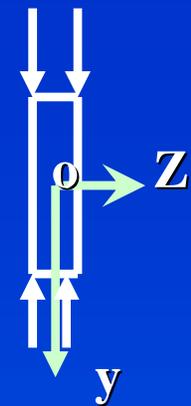
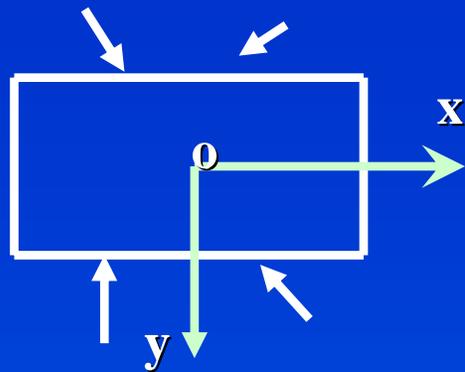
取中面为xy面，

在 $z = \pm t/2$ 处，横力 $T=0$



$$\sigma_z \Big|_{z=\pm t/2} = \tau_{zx} \Big|_{z=\pm t/2} = \tau_{zy} \Big|_{z=\pm t/2} = 0$$

由于连续性



→ $\sigma_z = \tau_{zx} = \tau_{zy} = 0$

同理: $\varepsilon, \gamma, u, v, w$ 不沿厚度变化

→ $\gamma_{zx} = \gamma_{zy} = 0$

$$\varepsilon_x, \varepsilon_y, \gamma_{xy}, u, v = ?$$

$$\varepsilon_z \neq 0, w \neq 0 \text{ 可求.}$$

沿一方向的尺寸很长,外力平行于截面且不沿长度变化

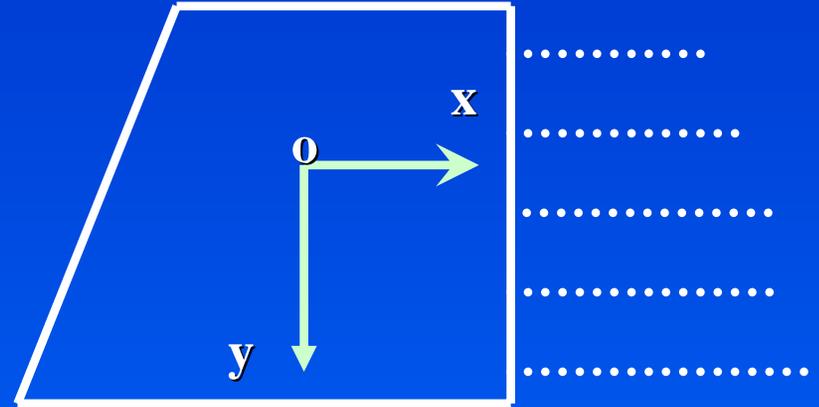
2. 平面应变 (平面位移)

$$w=0$$

由于对称性

$$\gamma_{zy} = \gamma_{zx} = 0$$

$$\varepsilon_z = 0, \tau_{zy} = \tau_{zx} = 0$$



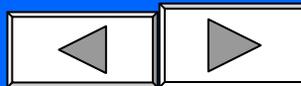
从三方面来分析:

静力学—研究物体内的平衡条件以及

静力边界条件

几何学—位移和应变的关系

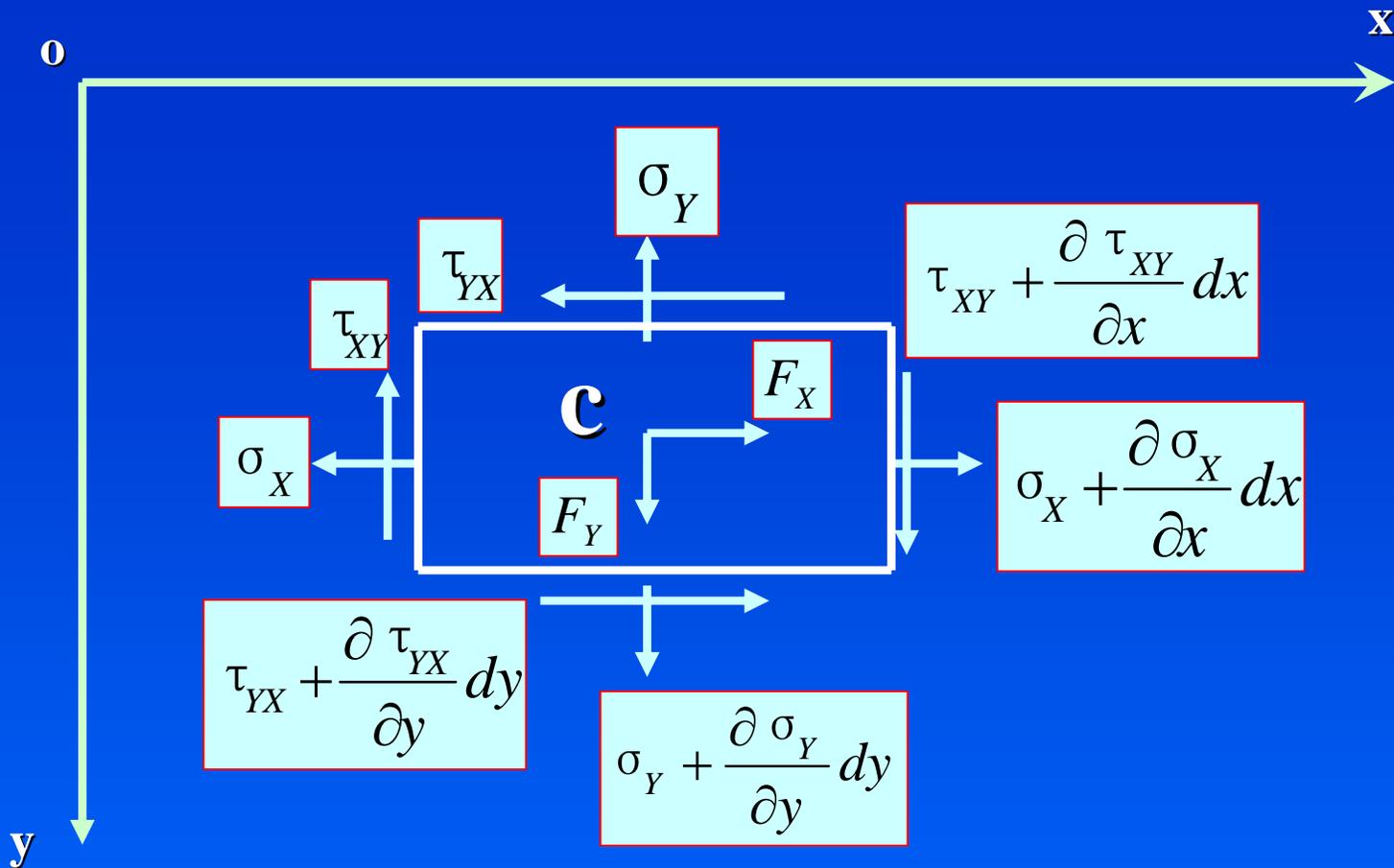
物理学—应力与应变的关系



§ 3.2 平衡微分方程

(Differential Equations of Equilibrium)





由 $\Sigma Mc=0$



$$\tau_{xy} = \tau_{yx}$$

由 $\Sigma X=0$

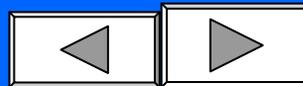


$$\frac{\partial \sigma_x}{\partial x} + \frac{\partial \tau_{yx}}{\partial y} + F_x = 0$$

由 $\Sigma Y=0$



$$\frac{\partial \tau_{xy}}{\partial x} + \frac{\partial \sigma_y}{\partial y} + F_y = 0$$



§ 3.3 一点的应力状态

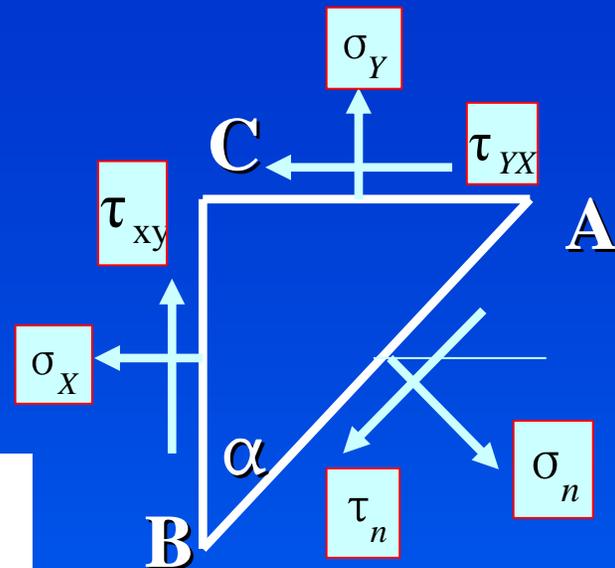
$$\begin{aligned} & \sigma_x ds \cos \alpha + \tau_{xy} ds \sin \alpha \\ & = \sigma_n ds \cos \alpha - \tau_n ds \sin \alpha \end{aligned}$$

$$\sigma_x l + \tau_{xy} m = \sigma_n l - \tau_n m \quad (1)$$

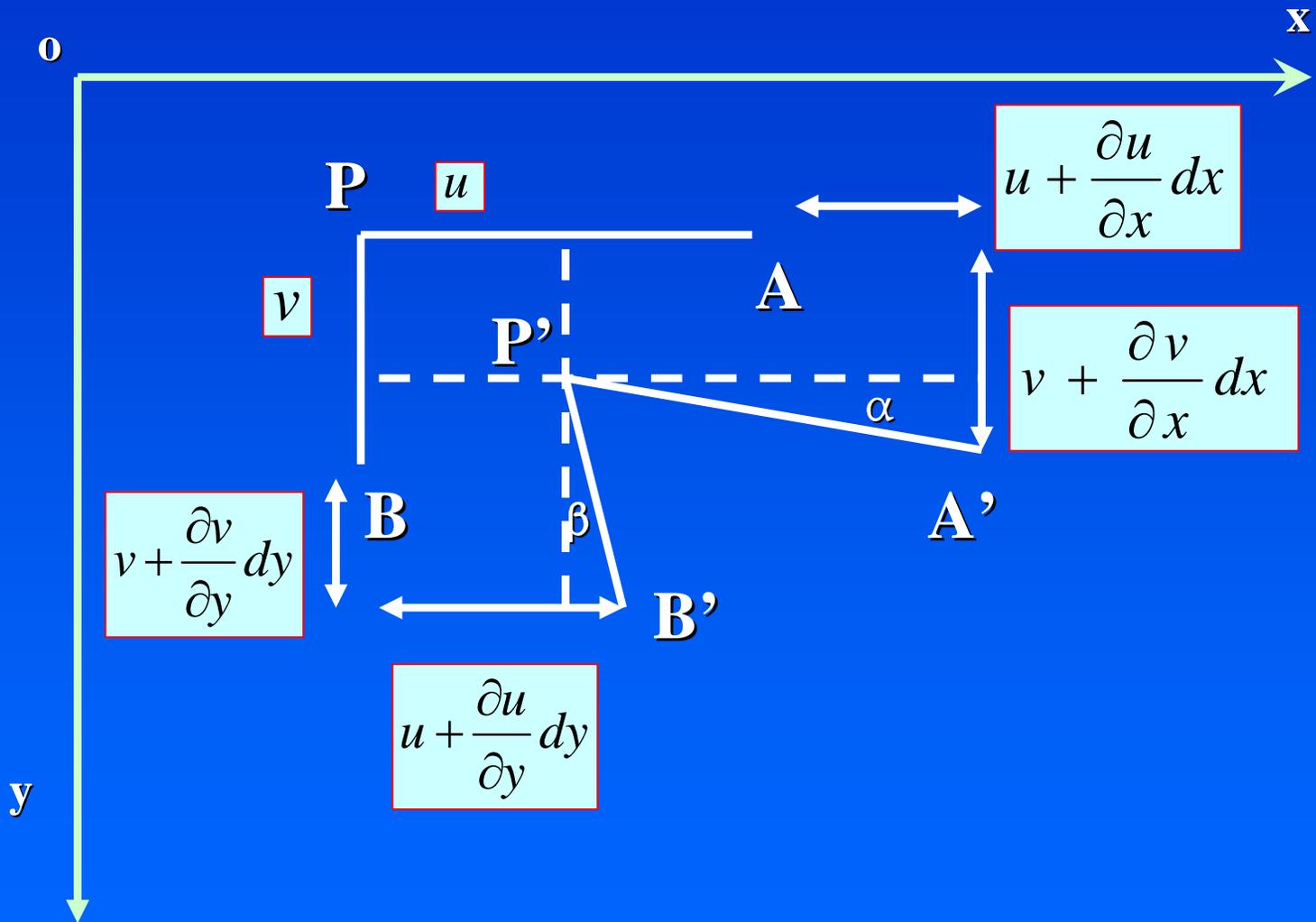
$$\tau_{xy} ds \cos \alpha + \sigma_y ds \sin \alpha$$

$$= \sigma_n ds \sin \alpha + \tau_n ds \cos \alpha$$

$$\tau_{xy} l + \sigma_y m = \sigma_n m + \tau_n l \quad (2)$$



§ 3.4 几何方程

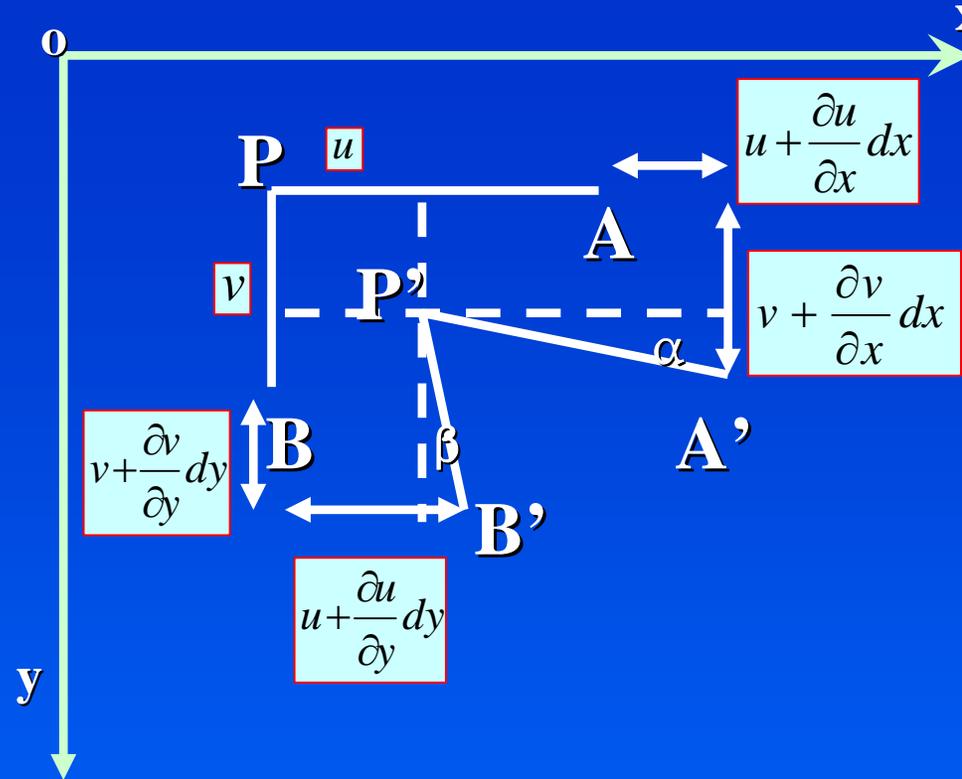


x方向正应变

$$\varepsilon_x = \frac{[u + (\partial u / \partial x) dx] - u}{dx} = \frac{\partial u}{\partial x}$$

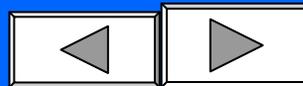
y方向正应变

$$\varepsilon_y = \frac{[v + (\partial v / \partial y) dy] - v}{dy} = \frac{\partial v}{\partial y}$$



剪应变 γ_{xy}

$$\gamma_{xy} = \alpha + \beta = \operatorname{tg} \alpha + \operatorname{tg} \beta = \frac{(v + \frac{\partial v}{\partial x} dx) - v}{dx} + \frac{(u + \frac{\partial u}{\partial y} dy) - u}{dy} = \frac{\partial v}{\partial x} + \frac{\partial u}{\partial y}$$



已知 u, v

$$\varepsilon_x, \varepsilon_y, \gamma_{xy},$$

$$\varepsilon_x = \frac{\partial u}{\partial x}$$

$$\varepsilon_y = \frac{\partial v}{\partial y}$$

$$\gamma_{xy} = \frac{\partial v}{\partial x} + \frac{\partial u}{\partial y}$$

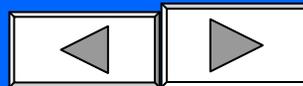
如

$$u = u_0 - \omega y, \quad v = v_0 + \omega x$$

$$\varepsilon_x = \varepsilon_y = \gamma_{xy} = 0$$

$$u_0, v_0, \omega$$

为刚体位移



变形协调方程或相容方程

(Compatibility Equation)

几何方程:

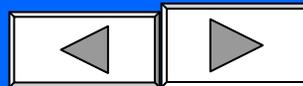
$$\varepsilon_x = \frac{\partial u}{\partial x}$$

$$\varepsilon_y = \frac{\partial v}{\partial y}$$

$$\gamma_{xy} = \frac{\partial v}{\partial x} + \frac{\partial u}{\partial y}$$

$$\frac{\partial^2 \varepsilon_x}{\partial y^2} + \frac{\partial^2 \varepsilon_y}{\partial x^2} = \frac{\partial^3 u}{\partial x \partial y^2} + \frac{\partial^3 v}{\partial x^2 \partial y} = \frac{\partial^2}{\partial x \partial y} \left(\frac{\partial u}{\partial y} + \frac{\partial v}{\partial x} \right) = \frac{\partial^2 \gamma_{xy}}{\partial x \partial y}$$

$$\frac{\partial^2 \varepsilon_x}{\partial y^2} + \frac{\partial^2 \varepsilon_y}{\partial x^2} = \frac{\partial^2 \gamma_{xy}}{\partial x \partial y}$$



§ 3.5 本构方程, 物理方程

应力--应变关系, HOOK'S定理

$$\varepsilon_x = \frac{1}{E} [\sigma_x - \mu (\sigma_y + \sigma_z)]$$

$$\varepsilon_y = \frac{1}{E} [\sigma_y - \mu (\sigma_x + \sigma_z)]$$

$$\varepsilon_z = \frac{1}{E} [\sigma_z - \mu (\sigma_x + \sigma_y)]$$

$$\gamma_{yz} = \frac{1}{G} \tau_{yz}$$

$$\gamma_{xz} = \frac{1}{G} \tau_{xz}$$

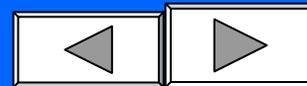
$$\gamma_{xy} = \frac{1}{G} \tau_{xy}$$

E--弹性模量

μ --泊松比

G--剪切模量

$$G = \frac{E}{2(1 + \mu)}$$



对平面应力问题

$$\sigma_z = \tau_{zx} = \tau_{zy} = 0$$

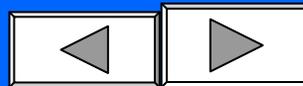
$$\varepsilon_x = (\sigma_x - \mu \sigma_y) / E$$

$$\varepsilon_y = (\sigma_y - \mu \sigma_x) / E$$

$$\gamma_{xy} = \tau_{xy} / G$$

$$\varepsilon_z = -\mu (\sigma_x + \sigma_y) / E$$

$$\gamma_{yz} = 0 \quad \gamma_{xz} = 0$$



平面应变问题

$$\varepsilon_z = \tau_{yz} = \tau_{zx} = 0 \quad \Rightarrow$$

$$\sigma_z = \mu (\sigma_x + \sigma_y), \quad \gamma_{yz} = \gamma_{zx} = 0$$

$$\begin{cases} \varepsilon_x = (1 - \mu^2) \left(\sigma_x - \frac{\mu}{1 - \mu} \sigma_y \right) / E \\ \varepsilon_y = (1 - \mu^2) \left(\sigma_y - \frac{\mu}{1 - \mu} \sigma_x \right) / E \\ \gamma_{xy} = \frac{2(1 + \mu)}{E} \tau_{xy} \end{cases}$$



平面应力问题

$$\varepsilon_x = (\sigma_x - \mu\sigma_y) / E$$

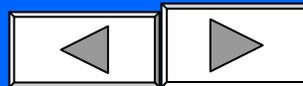
$$\varepsilon_y = (\sigma_y - \mu\sigma_x) / E$$

$$\gamma_{xy} = \tau_{xy} / G$$

平面应变问题

$$E \rightarrow \frac{E}{1 - \mu^2}, \quad \mu \rightarrow \frac{\mu}{1 - \mu}$$

$$\begin{cases} \varepsilon_x = (1 - \mu^2)(\sigma_x - \frac{\mu}{1 - \mu}\sigma_y) / E \\ \varepsilon_y = (1 - \mu^2)(\sigma_y - \frac{\mu}{1 - \mu}\sigma_x) / E \\ \gamma_{xy} = \frac{2(1 + \mu)}{E}\tau_{xy} \end{cases}$$



§ 3.6 边界条件, 圣维南原理

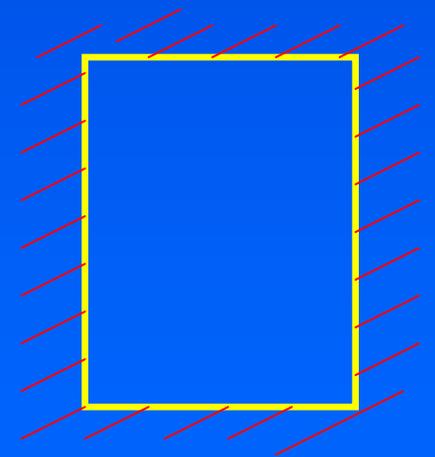
(Boundary Conditions,

Saint-Venant's Principle -1855)

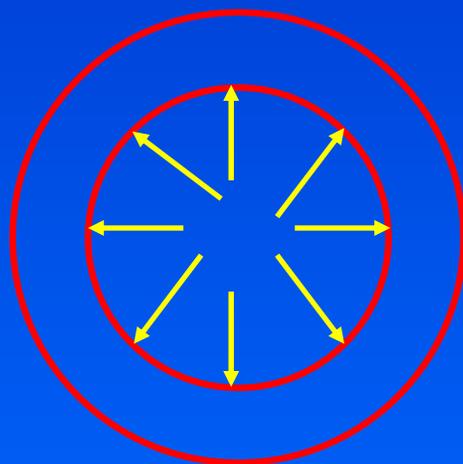
位移边界条件, 应力边界条件, 混合边界条件

1. 位移边界条件

$$\bar{u} = \bar{u}(\chi, \lambda) \quad \bar{v} = \bar{v}(\chi, \lambda)$$



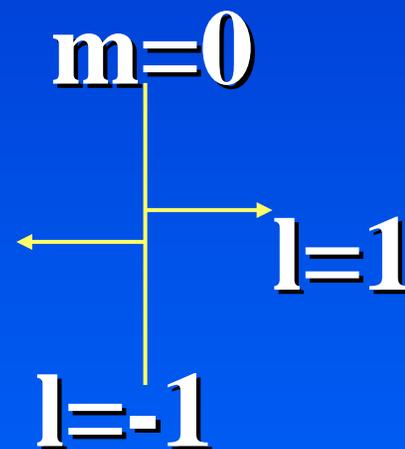
2.应力边界条件



$$\begin{cases} (\sigma_x l + \tau_{xy} m)_s = \overline{X}(x, y) \\ (\sigma_y m + \tau_{xy} l)_s = \overline{Y}(x, y) \end{cases}$$

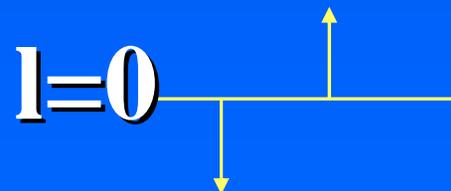
当 $l = \pm 1, m = 0$ 时

$$\begin{cases} (\sigma_x)_s = \pm \overline{X}(x, y) \\ (\tau_{xy})_s = \pm \overline{Y}(x, y) \end{cases}$$



$l = 0, m = \pm 1$

$$\begin{cases} (\tau_{xy})_s = \pm \overline{X}(x, y) \\ (\sigma_y)_s = \pm \overline{Y}(x, y) \end{cases}$$



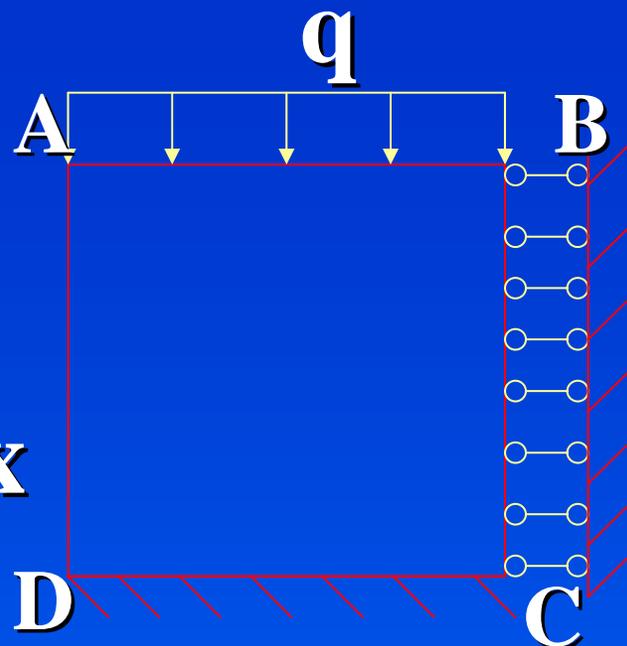
3.混合边界条件

应力边界 AB:

$$(\sigma_y)_s = -q, \quad (\tau_{xy})_s = 0$$

y

x



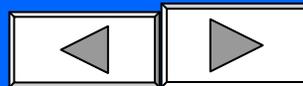
位移边界 CD:

$$(\bar{u})_s = (\bar{v})_s = 0$$

应力-位移边界 BC:

$$(\bar{u})_s = 0, \quad (\tau_{xy})_s = 0$$

弹性力学问题在数学上常被称为边界问题

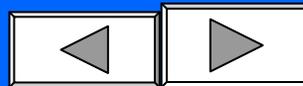


圣维南原

理。

两组静力等效，但分布不同的外力引起的应力场，仅在外力作用的附近有明显的不同，远处的应力场没有差别

静力等效：主矢量相同，对同一点的主矩也相同

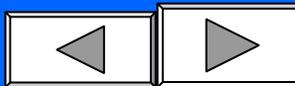


弹性力学问题的求解

平衡方程2个
几何方程3个
物理方程3个

应力分量: 3个
应变分量: 3个
位移分量: 2个

经典理论: 凑合解法
现代理论: 数值求解
(有限元法)



§ 3.6 虚功原理

虚功原理：在某一状态 $\{F\}$ ，应力 $\{\sigma\}$ 下，设有虚位移 $\{a^*\}$ ，对应的应变为 $\{\varepsilon^*\}$ ，则：

$$\iiint_{V_e} \{F\}^T \{a^*\} dV = \iiint_{V_e} \{\varepsilon^*\}^T \{\sigma\} dV$$

$$W_i = W_e$$

